



Article

Crystal Growth and Investigation of High-Pressure Physical Properties of Fe₂As

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Abstract: We reported the growth of Fe₂As single crystals and the study of its physical properties via comprehensive measurements, such as transport properties under pressure and high-pressure synchrotron radiation X-ray diffraction. Fe₂As is an antiferromagnetic metal with $T_N \sim 355$ K. Within the pressure range of 100 GPa, no superconductivity was observed above 2 K. The abrupt drop in resistance from 21 to 31.7 GPa suggests a high-pressure phase transition happens. The high-pressure X-ray experiments indicate a new high-pressure phase appears, starting from 27.13 GPa. After the refinement of the high-pressure X-ray data, the pressure dependence of lattice constants of Fe₂As (P4/nmm phase) was plotted and the bulk modulus B_0 was obtained to be 168.6 GPa.

Keywords: iron-based compound; high-pressure resistivity; high-pressure structure

1. Introduction

The discovery of the superconducting LaO_{1-x}F_xFeAs [1] with Tc = 26 K has attracted great interest to search for new iron-based superconductors. Soon after this discovery, a series of iron-based layered oxypnictides RFeAs(O_{1-x}F_x) were found by replacing La atoms with other rare-earth elements R (R = Ce, Pr, Nd, Sm, ...) [2–5]. The superconducting transition temperature was quickly enhanced to a maximum of 55 K in SmO_{1-x}F_xFeAs [4]. Subsequently, several families of iron-based superconductors were discovered, such as 122-type (Ba, K)Fe₂As₂[6], 11-type FeSe [7] and 111-type LiFeAs [8]. Iron-based superconductors contain the common Fe_{2×2} (X = As and Se) unit, which is believed to be crucial to support superconductivity in the iron-based materials. The superconductivity can be either induced by introducing charge carriers into the parent compounds or applying high pressure via suppressing the antiferromagnetic order [9–12]. Thus, the mechanism of superconductivity in iron-based compounds is considered to be highly related with the antiferromagnetic spin fluctuation.

Fe₂As crystallizes into a Cu₂Sb-type crystal structure. More interestingly, it contains [FeAs] layers, intercalated by Fe ions, and is isostructural to 111-type LiFeAs. For LiFeAs, there is no magnetic ordering and structural transition in compound LiFeAs, which makes it a special case in iron-based superconductors [13–15]. In contrast, Fe₂As demonstrates an antiferromagnetic transition with a Néel temperature of about 353 K [16]. Therefore, it is possible for Fe₂As to become a superconductor if the antiferromagnetic order is suppressed. High pressure was applied to Fe₂As to

Crystals **2020**, 10, 790

explore the superconductivity within 32 GPa. However, no sign of superconductivity was observed [17,18]. In this work, we report on the study of properties of Fe₂As within 100 GPa. Fe₂As keeps its metallic state and no superconductivity is observed above 2 K. It undergoes a high-pressure phase transition at about 27 GPa, which rules out the possibility of pressure-induced superconductivity in the tetragonal P4/nmm phase of Fe₂As. Via the study of high-pressure structure, the bulk modulus B_0 for the P4/nmm phase of Fe₂As is obtained to be about 168.6 GPa.

2. Materials and Methods

Single crystals of Fe₂As were grown using the self-flux method. High-purity Fe powder (Alfa, 99.999%) and As powder (Alfa, 99.999%) were mixed with a molar ratio of 2:1. Then the mixtures were put in an Al₂O₃ capsule, sealed into a niobium tube filled with high-purity Ar gas and then sealed in an evacuated quartz tube. The quartz tube was heated to 1180 °C at a rate of 4 °C/min and hold at this temperature for 20 h. Then, it was slowly cooled down to 600 °C at a rate of 3 °C/h, after which the single crystals of Fe₂As with the typical size of 2 mm were obtained.

The samples were characterized by x-ray powder diffraction (XRD) in the 10– 90° range on a Philips X'pert diffractometer (PANalytical B.V., Almelo, Netherlands) using Cu K α radiation generated at 40 kV and 40 mA. The chemical compositions were examined by energy dispersive X-ray spectroscopy (EDX). The DC magnetic susceptibility was measured using a superconducting quantum interference device (SQUID) (Quantum Design, San Diego, CA, USA). The electrical resistivity under ambient pressure was measured by a physical property measurement system (PPMS) (Quantum Design, San Diego, CA, USA) using the four-probe method. The four Pt lines were pasted using silver conductive adhesive onto the flat surface (ab-plane) of the single crystal sample, thus we could get the ab-plane resistance when current was applied in the ab-plane.

The electronic transport properties under high pressure were investigated using the four-probe electrical conductivity method as described in [19,20]. The dimension of the sample used in our measurement was about 35 μ m \times 35 μ m for the surface (*ab*-plane) of single crystal and 5 μ m thick, so that the four electrodes could be contacted tightly with the flat surface. NaCl powder was used as the pressure transmitting medium. The pressure was calibrated via the ruby fluorescence method at room temperature [21]. After being fully assembled, we put the diamond anvil cell inside a MagLab system (Oxford Instruments, Oxford, UK) to perform the electric transport experiments within the range of 2 K to 300 K.

In situ high pressure synchrotron X-ray diffraction experiments were performed at Beijing Synchrotron Radiation Facility (Beijing, China) with a wavelength of 0.6199 Å using a symmetric Mao Bell DAC (MHD Corporation, Fukuoka, Japan). Fine grounded polycrystalline powders were loaded in DAC and ruby balls were placed around the sample as pressure markers. The obtained two-dimensional image plate patterns were converted to one-dimensional 20 versus intensity data using the Fit2D software package (Andy Hammersley, ESRF, Grenoble, France) [22]. Refinements of the high-pressure synchrotron X-ray diffraction patterns were performed by Rietveld method [23] through the GSAS package (Lansce Los Alamos, Los Alamos, NM, USA).

3. Results and Discussion

The compound Fe₂As crystallizes into the Cu₂Sb-type crystal structure. Figure 1a shows the X-ray diffraction (XRD) pattern of Fe₂As single crystal. Only sharp [00*l*] peaks can be observed, suggesting a good quality of our crystal sample. A Le Bail fit to the diffraction data yields c = 5.965 Å, which coincides with the previous result [24] and is much smaller than that of LiFeAs (c = 6.3574 Å) [8]. The photograph of Fe₂As crystal with the size up to 1 mm × 2 mm × 0.5 mm is shown in the inset of Figure 1a. The elemental composition of the single crystal was checked by EDX, giving an average atomic ratio of Fe:As = 2.07:1, which is very close to the elemental ratio of stoichiometric Fe₂As. Figure 1b shows the DC magnetic susceptibility measurements on Fe₂As single crystal from 5 K to 380 K with applied magnetic field parallel to ab-plane in both zero-field cooled (ZFC) and field cooled (FC) modes. We can unambiguously find an antiferromagnetic transition at about 355 K, which is in agreement with the result of magnetic susceptibility and neutron diffraction

Crystals **2020**, 10, 790 3 of 7

measurements [16]. Figure 1c shows the electrical resistivity versus temperature for Fe₂As single crystal at ambient pressure. An anomaly at about 355 K can be observed in the resistivity curve, which can be attributed to the antiferromagnetic transition.

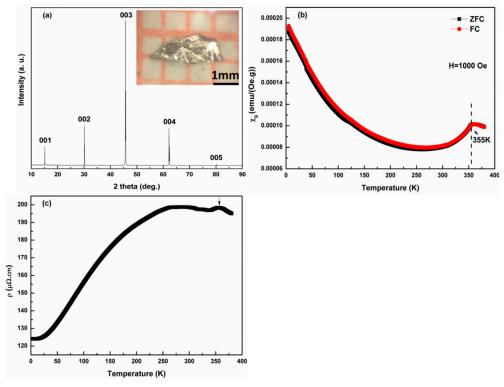


Figure 1. (a) X-ray diffraction (XRD) pattern of Fe₂As single crystal. The inset shows the photograph of Fe₂As single crystal; (b) Zero-field cooled (ZFC) and field cooled (FC) susceptibilities versus temperature for Fe₂As single crystal with applied magnetic field parallel to ab-plane, an obvious antiferromagnetic transition can be seen at about 355 K; (c) Temperature dependence of *ab* plane resistance at ambient pressure. An obvious anomaly at about 355 K in the resistivity curve was marked by arrow.

For the iron-based superconductors, superconductivity can always be induced by suppressing antiferromagnetic order. Pressure is a clean tool to suppress the antiferromagnetic order without introducing defects or impurities and, then, induce the possible superconductivity. The research on tuning superconductivity by pressure has been widely reported in many material systems ranging from topological materials to iron-based superconductors [25-33]. Therefore, we conducted the measurement of resistance for Fe₂As within the range of 2 K to 300 K, under pressure as high as 100 GPa, which can be seen in Figure 2. As reported before [17,18], it is difficult to determine the Néel temperature from the high-pressure resistance curve. Within the pressure range of 21 GPa, the resistance at low temperature increases as pressure increases, which should arise from the suppression of antiferromagnetic order and the enhancement of magnetic fluctuation. While within lower pressure the decrease in resistance with increasing pressure was reported in the references [17,18]. The disagreement is suggested to be caused by the difference in samples. In our work, single crystal was used for the resistance measurement and the data we got indicate ab-plane resistance, while the polycrystalline samples were used in the previous work where the extrinsic grains contacting resistance is included. When pressure increases from 21 to 31.7 GPa, the resistance drops abruptly. Generally, high-pressure phase has a more compacted structure and leads to a lower resistance. In addition, the new phase might completely destroy the antiferromagnetic order and push Fe₂As away from antiferromagnetic fluctuation, which can also cause the drop in resistance. Therefore, it is hinted that the sharp drop in resistance is associated with the occurrence of pressure-induced structure transition as will be confirmed by the following high-pressure X-ray experiments. Above 31.7 GPa, the resistance gradually increases when further increasing the

Crystals **2020**, 10, 790 4 of 7

pressure. For Fe₂As, the anomaly of pressure dependence of resistance has also been observed in the pressure range from 10.4 to 31.4 GPa in the previous work [17]. In any case, we cannot observe any sign of superconductivity above 2 K within our highest experimental pressure.

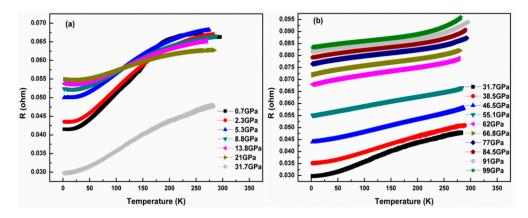


Figure 2. (a) Temperature dependence of *ab*-plane resistance at various pressures range from 0.7 to 31.7 GPa; (b) Temperature dependence of *ab*-plane resistance at various pressures range from 31.7 to 99 GPa.

To examine the high-pressure phase transition, we carried out the high-pressure synchrotron radiation X-ray experiments. Figure 3 shows the X-ray diffraction patterns under high pressure. Below 23.5 GPa, all the peaks can be indexed very well as a tetragonal *P*4/nmm phase. When the pressure increases up to 27.13 GPa, several peaks derived from *P*4/nmm phase become weak and disappear gradually. Simultaneously, there appears a set of new diffraction peaks marked with inverted triangles, as can be seen in Figure 3b, which indicates the appearance of the high-pressure phase. Further increasing the pressure, the high-pressure phase gradually becomes the main phase.

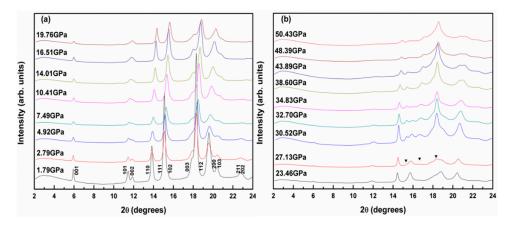


Figure 3. Synchrotron X-ray diffraction patterns of Fe₂As polycrystalline at different pressures up to 50.43 GPa (λ = 0.6199 Å). (**a**) XRD patterns at various pressures range from 1.79 to 19.76 GPa; (**b**) XRD patterns at various pressures range from 23.46 to 50.43 GPa, the peaks marked with inverted triangles denote the new diffraction peaks emerging.

Data analysis of the diffraction patterns was performed with the Rietveld refinements using the GSAS program package, from which we can obtain the values of lattice constants and unit cell volumes under different pressures for P4/nmm phase. As an example, Figure 4 presents the Rietveld refinements for the diffraction patterns at 1.79 GPa. The evolution of the lattice parameters with a pressure for P4/nmm phase is shown in Figure 5a. Both the lattice constants of a and c are compressed by pressure. Within 25 GPa, a and c shrink about 3.9% and 1.6%, respectively; this demonstrates that the ab-plane is easier to be compressed than c axis. Figure 5b presents the unit cell

Crystals **2020**, 10, 790 5 of 7

volumes versus pressure. The volume–pressure data are fitted using the third-order Birch–Murnaghan equation of states, $P(GPa) = \frac{3}{2} \times B_0 \left[\left(\frac{V_0}{V} \right)^{\frac{7}{3}} - \left(\frac{V_0}{V} \right)^{\frac{5}{3}} \right] \times \left\{ 1 - \left(3 - \frac{3}{4} \times B_0' \right) \times \left[\left(\frac{V_0}{V} \right)^{\frac{2}{3}} - 1 \right] \right\}$

when the first pressure derivative B'_0 is fixed to be 4, the fitting gives the value of bulk modulus $B_0 = 168.6$ GPa, which is higher than LiFeAs and close to that of CuFeAs [34].

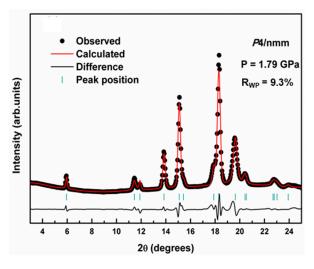


Figure 4. Rietveld refined XRD patterns of Fe₂As at 1.79 GPa.

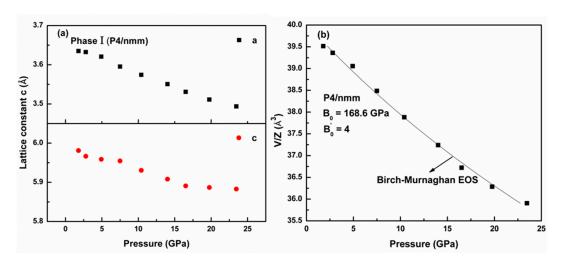


Figure 5. (a) Pressure dependence of the lattice constants *a* and *c* for *P*4/nmm phase; (b) Unit cell volumes versus pressure. The solid line shows the fitting curve using the third-order Birch–Murnaghan equation of states.

For the unknown high-pressure phase of Fe₂As, we can refer to the previous works about high-pressure phase transition for the isostructural "111" type iron based superconductors. NaFeAs has been reported to undergo an isostructural phase transition at 20 Gpa [35]. Interestingly, the high-pressure crystal structures prediction of "111" type iron-based superconductors (LiFeAs, NaFeAs, and LiFeP) based on CALYPSO (Crystal structure Analysis by Particle Swarm Optimization) methodology indicates there exist two possible high-pressure crystal structures below 100 GPa, the Cmcm structure for NaFeAs or LiFeP and the *P*-3m1 structure for LiFeAs [36]. However, using either Cmcm or *P*-3m1 structure as the initial model to carry out the refinements for our high-pressure X-ray data, we found that the peaks of high-pressure phase cannot be well fitted. Thus, the high-pressure phase of Fe₂As needs further studies to be determined. In any case, the high-pressure phase transition rules out the possibility of pressure-induced superconductivity in the tetragonal *P*4/nmm phase of Fe₂As.

Crystals **2020**, 10, 790 6 of 7

4. Conclusions

In this work, high quality single crystals of Fe₂As were grown and the properties under high pressure were studied. Within 100 GPa, Fe₂As keeps its metallic state and no superconductivity is observed above 2 K. A high-pressure phase transition happens at about 27 GPa, which suggests that pressure-induced superconductivity cannot be observed in the tetragonal P4/nmm phase of Fe₂As. At last, via the study of structure under high pressure we obtained the bulk modulus B_0 of about 168.6 GPa for the ambient phase of Fe₂As.

Author Contributions: X.W. and C.J. conceived and designed the experiments and edited the manuscript; G.D. performed the experiments, analyzed the data and drafted the manuscript; C.Z. and S.Z. contributed the initial high-pressure resistivity measurements; X.H., C.R., J.Z. and X.L. contributed the high-pressure synchrotron radiation X-ray experiments. All authors have read and agreed to the published version of the manuscript.

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