

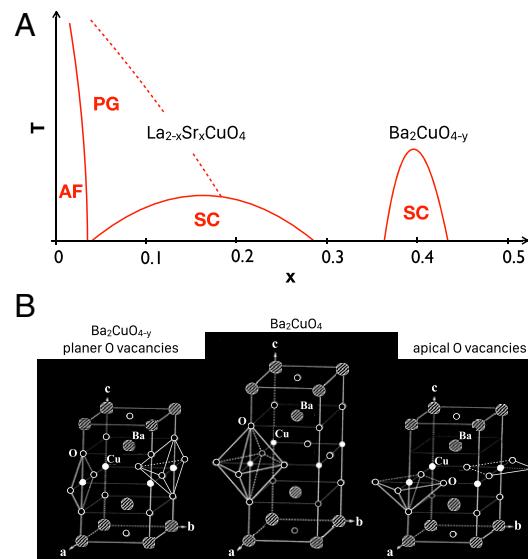
# A different branch of the high $T_c$ family?

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In PNAS, Li et al. (1) suggest that  $\text{Ba}_2\text{CuO}_{4-y}$  is a member of a different branch of high- $T_c$  cuprate superconducting materials. This branch is characterized as heavily overdoped with an exceptionally short Cu apical O spacing and O vacancies that are located in the  $\text{CuO}_2$  planes. These characteristics, illustrated in Fig. 1, differ in significant ways from those of the traditional cuprate superconducting materials (2).

For example, consider  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  (LSCO), which has the same  $\text{K}_2\text{NiF}_4$  structure as  $\text{Ba}_2\text{CuO}_{4-y}$ . Using a valence count with  $\text{La}^{+2}$ ,  $\text{Sr}^{+3}$ , and  $\text{O}^{-2}$ , the number of holes in the Cu 3d shell of LSCO is  $1+x$ . For  $x=0$ , LCO is antiferromagnetic. As Sr is added the antiferromagnetic Néel temperature decreases as shown in Fig. 1A and superconductivity onsets with  $T_c$  peaking at a small doping  $x \sim 0.15$ . At larger doping, as the system moves further away from the antiferromagnetic Mott-Hubbard phase,  $T_c$  decreases and vanishes for  $x > 0.25$ . For  $\text{Ba}_2\text{CuO}_{4-y}$ , the valence of Ba is +2 so that the hole doping  $x = 2(1-y)$ . Thus, the  $y=0.8$  sample Li et al. (1) discuss is heavily overdoped with  $x=0.4$ . Nevertheless, the  $T_c$  of  $\text{Ba}_2\text{CuO}_{4-y}$  is approximately 30 K higher than that of optimally doped LSCO.

Like the previously reported highly overdoped cuprate materials (3–8)  $\text{Sr}_2\text{CuO}_{4-y}$ ,  $(\text{Sr},\text{Ba})_2\text{CuO}_{4-y}$ , and  $\text{Cu}_{0.75}\text{Mo}_{0.25}\text{Sr}_2\text{YCu}_2\text{O}_{7.54}$ ,  $\text{Ba}_2\text{CuO}_{4-y}$  is synthesized under high pressure and high temperature in the presence of a strong oxidizing agent. While the crystal has the  $\text{K}_2\text{NiF}_4$  structure illustrated in the center of Fig. 1B, the  $\text{CuO}_6$  octahedron structure is highly compressed with a much shorter Cu-O apical distance (1.86 Å) than the traditional cuprates (2.42 Å for  $\text{La}_2\text{CuO}_4$ ). This is schematically illustrated at the left and right in Fig. 1B. The shortening of the Cu-O apical distance raises the energy of the  $3d_{3z^2-r^2}$  orbital so that in  $\text{Ba}_2\text{CuO}_{4-y}$  the Cu states near the Fermi level have both  $3d_{3z^2-r^2}$  and  $3d_{x^2-y^2}$  orbital character. As the authors note, a shortened apical Cu-O distance and the admixing of  $3d_{3z^2-r^2}$  orbital weight in the states near the Fermi energy are found to reduce  $T_c$  in the traditional cuprate superconductors (9).



**Fig. 1.** (A) A schematic cuprate phase diagram as a function of hole doping  $x$ . On the left at low doping are the familiar antiferromagnetic (AF), pseudogap (PG), and  $d$ -wave superconducting (SC) regions of  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ . On the right is the heavily overdoped  $x \sim 0.4$  region of superconductivity Li et al. (1) discuss for  $\text{Ba}_2\text{CuO}_{4-y}$ . (B) The center panel shows a fictitious  $\text{Ba}_2\text{CuO}_4$  compound with the same  $\text{K}_2\text{NiF}_4$  structure as  $\text{La}_2\text{CuO}_4$ . The structures shown on the left and right have the compressed  $c$ -axis structure of  $\text{Ba}_2\text{CuO}_{4-y}$ . In this case the hole doping is controlled by the O vacancies. For the structure on the right, the O vacancies are on the apical sites and the  $\text{CuO}_4$  corner-shared  $\text{CuO}_2$  sheets remain intact. However, according to Li et al. (1) in  $\text{Ba}_2\text{CuO}_{4-y}$  the O vacancies are in the plane. In this case, as shown by the structure on the left, the  $\text{CuO}_2$  sheets are destroyed and randomly oriented Cu-O chain segments are likely formed.

Finally, the authors note at the end of the legend for figure 3 in ref. 1 that while the exact positions of the O vacancies are not known at present, they are in the  $\text{CuO}_2$  planes. This agrees with the conclusions of Geballe and Marezio (5) regarding the O vacancies in  $\text{Sr}_2\text{CuO}_{4-y}$ . A number of earlier studies Li et al. (1) cite assume that the oxygen vacancies were at the apical O sites. In this

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case, while there would be CuO<sub>5</sub> pyramidal structures and square fourfold coordinated CuO<sub>4</sub> along with the CuO<sub>6</sub> octahedral units, the CuO<sub>2</sub> planes would survive intact as illustrated in the right-hand structure of Fig. 1B. However, if the oxygen vacancies are in the CuO<sub>2</sub> planes one will have Cu–O chain structures rather than the CuO<sub>2</sub> planes as indicated in the left-hand structure of Fig. 1B. The 2D CuO<sub>2</sub> planes, consisting of corner-shared CuO<sub>4</sub> units, are considered the key structural elements of the traditional cuprate superconductors. They are the defining characteristic of these superconductors and imperfections in these layers are known to reduce  $T_c$ .

As noted, Ba<sub>2</sub>CuO<sub>4-y</sub> is made under high pressure and temperature resulting in polycrystalline samples and the precise

location of the oxygen vacancies in the CuO<sub>2</sub> plane remains open. Nevertheless, the remarkably high  $T_c$  of this highly overdoped cuprate with its short Cu apical O separation and its O vacancies in the CuO<sub>2</sub> plane suggest that it is a member of a different branch of high- $T_c$  cuprate materials, which challenges the basic tenants of many high- $T_c$  theories.

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- 1 W. M. Li *et al.*, Superconductivity in a unique type of copper oxide. *Proc. Natl. Acad. Sci. U.S.A.* **116**, 12156–12160 (2019).
- 2 S. K. Gupta, H. Jangam, N. Sharma, Review of high temperature superconductors and application in various fields. *Res. Dev. Mater Sci.* **7**, 000668 (2018).
- 3 Z. Hiroi, M. Takano, M. Azuma, Y. Takeda, A new family of copper oxide superconductors Sr<sub>n+1</sub>Cu<sub>n</sub>O<sub>2n+1+δ</sub> stabilized at high pressure. *Nature* **364**, 315–317 (1993).
- 4 Q. Q. Liu *et al.*, Enhancement of the superconducting critical temperature of Sr<sub>2</sub>CuO<sub>3+δ</sub> up to 95 K by ordering dopant atoms. *Phys Rev B* **74**, 100506(R) (2006).
- 5 T. H. Geballe, M. Marezio, Enhanced superconductivity in Sr<sub>2</sub>CuO<sub>4-y</sub>. *Physica C* **469**, 680–684 (2009).
- 6 W. B. Gao *et al.*, Out-of-plane effect on the superconductivity of Sr<sub>2-x</sub>Ba<sub>x</sub>CuO<sub>3+δ</sub> with T<sub>c</sub> up to 98 K. *Phys. Rev. B* **80**, 094523 (2009).
- 7 Y. Liu *et al.*, A new modulated structure in Sr<sub>2</sub>CuO<sub>3+δ</sub> superconductor synthesized under high pressure. *Physica C* **497**:34–37 (2014).
- 8 A. Gauzzi *et al.*, Bulk superconductivity at 84 K in the strongly overdoped regime of cuprates. *Phys. Rev. B* **94**, 180509 (R) (2016).
- 9 E. Pavarini *et al.*, Band-structure trend in hole-doped cuprates and correlations with T<sub>cmax</sub>. *Phys. Rev. Lett.* **87**, 047003 (2001).