



# Structural stability of $\text{Zn}_3\text{N}_2$ under high pressure

J.G. Zhao<sup>a,b,\*</sup>, L.X. Yang<sup>b</sup>, S.J. You<sup>b</sup>, F.Y. Li<sup>b</sup>, C.Q. Jin<sup>b</sup>, J. Liu<sup>c</sup>

<sup>a</sup> Natural Science Research Center, Academy of Fundamental and Interdisciplinary Sciences, Harbin Institute of Technology, Harbin 150080, PR China

<sup>b</sup> Beijing National Lab for Condensed Matter Physics, Institute of Physics, Chinese Academy of Sciences, Beijing 100190, PR China

<sup>c</sup> Institute of High Energy Physics, Chinese Academy of Sciences, Beijing 100039, PR China

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## ABSTRACT

*In situ* high-pressure energy dispersive X-ray diffraction experiment on  $\text{Zn}_3\text{N}_2$  has been performed by using a diamond anvil cell instrument with synchrotron radiation at room temperature. The results showed that the structure of  $\text{Zn}_3\text{N}_2$  is stable in the experimental pressure range up to 25.2 GPa. According to the Birch–Murnaghan equation of state determined from the relationship of unit cell volume and pressure, assuming  $B'_0=4$ , the bulk modulus  $B_0=228(2)$  GPa has been obtained.

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## 1. Introduction

$\text{Zn}_3\text{N}_2$ , which is black in color, was first synthesized by Juza and Hahn in 1940 [1]. In 1993, Kuriyama et al. obtained the polycrystalline  $\text{Zn}_3\text{N}_2$  films through the direct reaction between Zn and  $\text{NH}_3$  which were evaporated onto quartz substrates [2]. At room temperature,  $\text{Zn}_3\text{N}_2$  film shows a high electron mobility of about  $100\text{ cm}^2/(\text{Vs})$  [3]. It was determined to be an n-type semiconductor, with a direct gap of 1.23 eV.  $\text{Zn}_3\text{N}_2$  film can be used to produce p-type  $\text{ZnO:N}$  [4]. Due to the possible application prospect,  $\text{Zn}_3\text{N}_2$  has attracted more and more attention since 2000. Recently, several groups obtained  $\text{Zn}_3\text{N}_2$  films through different methods, including shielded reactive vacuum arc deposition [5,6], molecular-beam epitaxy [7], reactive radio-frequency (rf) magnetron sputtering [8–11], etc.

$\text{Zn}_3\text{N}_2$  is isostructural with  $\text{Mg}_3\text{N}_2$ ,  $\text{Ca}_3\text{N}_2$ , and  $\text{Cd}_3\text{N}_2$ , which adopt the anti-bixbyite structure [12]. The space group is  $Ia\bar{3}$ , and lattice parameter  $a$  is equal to  $9.7691(1)\text{ \AA}$ .  $\text{Zn}_3\text{N}_2$  is a derivative of the  $\text{CaF}_2$  structure, where N and Zn atoms occupy the Ca positions and three-fourth of the F positions, respectively. In this cubic structure, Zn atoms occupy the tetrahedral sites of an approximately cubic close packed array of N atoms. The close-grained structure indicates the possible small compress coefficient of  $\text{Zn}_3\text{N}_2$ . Generally, the lattice mismatch as well as the different thermal expansion coefficients between thin film and substrate always lead to strain. Many nitrides undergo the crystal structure

phase transition under high pressure, such as  $\text{Li}_3\text{N}$ ,  $\text{Cu}_3\text{N}$ ,  $\text{Mg}_3\text{N}_2$ ,  $\text{Al}_3\text{N}$  [13–16]. Considering the application prospect of film materials, it is important to learn the pressure effect on the structure and consequently physical behaviors of  $\text{Zn}_3\text{N}_2$ . In this paper, we report the study of high-pressure structural stability of  $\text{Zn}_3\text{N}_2$  based on the diamond anvil cell (DAC) technique by using *in situ* high-pressure energy dispersive X-ray diffraction with synchrotron radiation.

## 2. Experiments

The *in situ* high-pressure X-ray energy dispersive diffraction experiment on  $\text{Zn}_3\text{N}_2$  was carried out at room temperature in a diamond anvil cell at Beijing Synchrotron Radiation Facility (BSRF). The culet of diamond is  $500\text{ }\mu\text{m}$  in diameter and the hole in a T301 stainless steel gasket is  $250\text{ }\mu\text{m}$  in diameter. The powder sample is loaded into the hole in the gasket. The pressure on the sample was measured by using the ruby fluorescence method. The spot size of the focused X-ray beam was  $25\times 30\text{ }\mu\text{m}$  and the storage ring was operated at 2.2 GeV and 60–100 mA. In this experiment, the relation of energy of photon and channel was  $E=0.5444+0.00885\text{ chn}$ . The diffraction angle  $\theta$  is fixed to  $10^\circ$ .

## 3. Results and discussions

Fig. 1 shows the spectra of X-ray energy dispersive diffraction results of  $\text{Zn}_3\text{N}_2$  under various pressures. The diffraction data are collected in the pressure range of 0–25.2 GPa. There are the diffraction peaks of  $\text{Zn}_3\text{N}_2$  and fluorescence peaks of Zn in the X-ray diffraction patterns. There also is an escaping peak

\* Corresponding author at: Natural Science Research Center, Academy of Fundamental and Interdisciplinary Sciences, Harbin Institute of Technology, Harbin 150080, PR China. Tel./fax: +86 451 86403249.

E-mail address: [zhaojingeng@163.com](mailto:zhaojingeng@163.com) (J.G. Zhao).

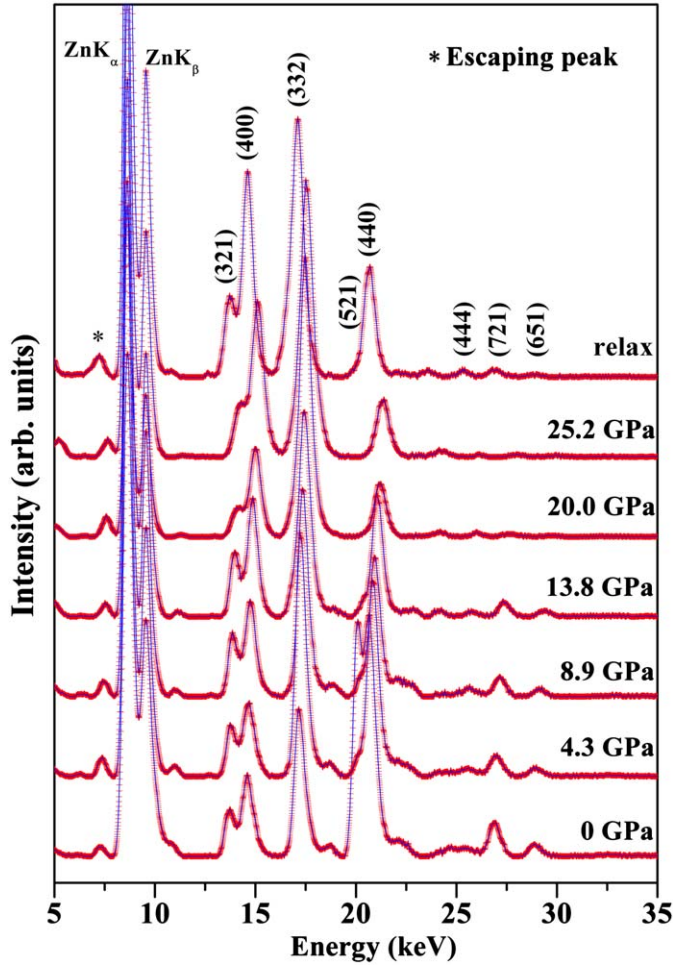


Fig. 1. Spectra of energy dispersive X-ray diffraction pattern of  $\text{Zn}_3\text{N}_2$  under various pressures.

indicated with asterisk, due to the high rank of the peak (3 3 2) of the sample. The indexes of sample peaks are denoted in the XRD patterns. All the Gaussian-type diffraction peaks shift towards the high-energy direction with increasing pressure and return to the original sites when the pressure releases to ambient. From the XRD patterns, the structural phase transition cannot be found in this experimental pressure range because no splitting or merging peaks are observed. Thus the peak fitting, indexing, and cell parameter refining are based on the primal structure of  $\text{Zn}_3\text{N}_2$  with the space group  $Ia-3$ .

According to the Bragg formula

$$E(\text{keV}) \times d(\text{\AA}) = 6.19925 / \sin \theta, \quad (1)$$

we got the  $d$  spacings of  $\text{Zn}_3\text{N}_2$  under various pressures. The lattice parameter and unit cell volume ( $V$ ) under different pressures are obtained through the  $d$  spacings. Fig. 2 shows the pressure dependences of unit cell volume and lattice parameter for  $\text{Zn}_3\text{N}_2$ , where the bold line is the fit to data. In the experimental error range,  $V$  is decreasing with the increasing pressure. Using the Birch–Murnaghan equation of state (EOS) [17]

$$P(\text{GPa}) = \frac{3}{2} \times B_0 \times \left[ \left( \frac{V}{V_0} \right)^{-7/3} - \left( \frac{V}{V_0} \right)^{-5/3} \right] \times \left\{ 1 - \left( 3 - \frac{3}{4} \times B'_0 \right) \times \left[ \left( \frac{V}{V_0} \right)^{-2/3} - 1 \right] \right\} \quad (2)$$

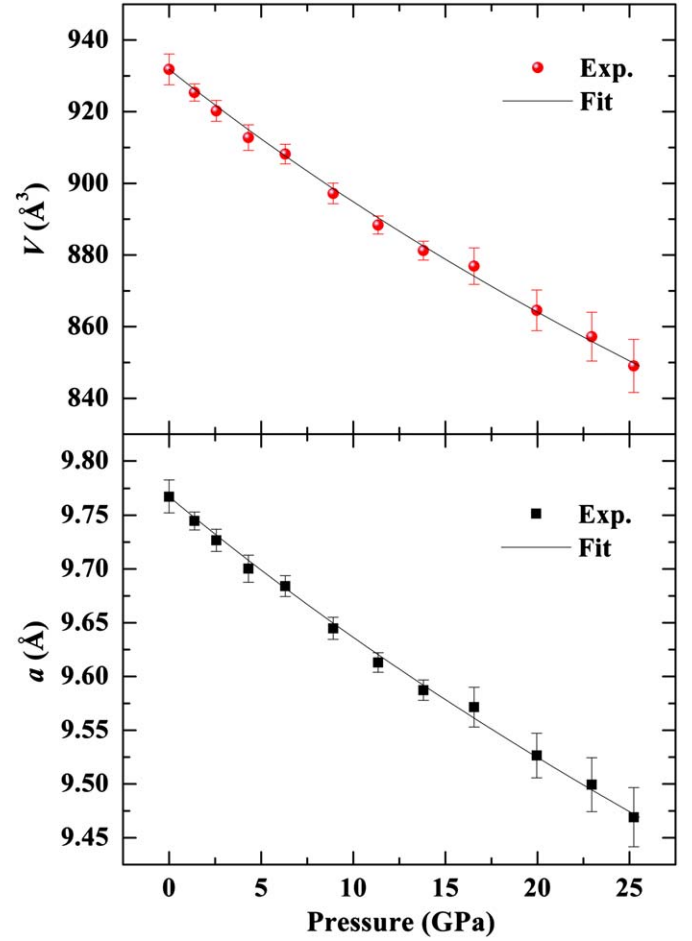


Fig. 2. Pressure dependences of (a) unit cell volume and (b) lattice parameter of  $\text{Zn}_3\text{N}_2$  at room temperature.

we obtained, assuming the first-order derivative to be  $B'_0=4$ , the ambient pressure bulk modulus  $B_0=228(2)$  GPa for  $\text{Zn}_3\text{N}_2$ . The value of  $B_0$  is larger than those of most nitrides, so  $\text{Zn}_3\text{N}_2$  is a stiff material.

The schematic view of  $\text{Zn}_3\text{N}_2$ , with the anti-bixbyite structure, is shown in Fig. 3(a). There are four N atoms around one Zn atom, which forms a  $\text{ZnN}_4$  tetrahedron. For the N atom, the surrounding six Zn atoms combine to one “ $\text{NZn}_6$  octahedron”, which is not a strict octahedron due to the smaller ion radius of Zn ion than that of N ion. The details of  $\text{ZnN}_4$  tetrahedron and “ $\text{NZn}_6$  octahedron” are shown in Fig. 3(b). According to Shannon Table [18], the ratio of the radius of  $\text{Zn}^{2+}$  ion to that of  $\text{N}^{3-}$  ion is equal to 0.411, being close to the ratio limit of 0.414 to form tetrahedron. So it is difficult to compress the  $\text{ZnN}_4$  tetrahedron, which induces that  $\text{Zn}_3\text{N}_2$  has a large ambient pressure bulk modulus.

#### 4. Conclusions

In summary, the structural stability of  $\text{Zn}_3\text{N}_2$  under high pressure has been investigated through *in situ* high-pressure X-ray energy dispersive diffraction experiments. There is no structural transition in the experimental pressure range. The big bulk modulus  $B_0$  of 228(2) GPa indicates that  $\text{Zn}_3\text{N}_2$  is a stiff nitride.

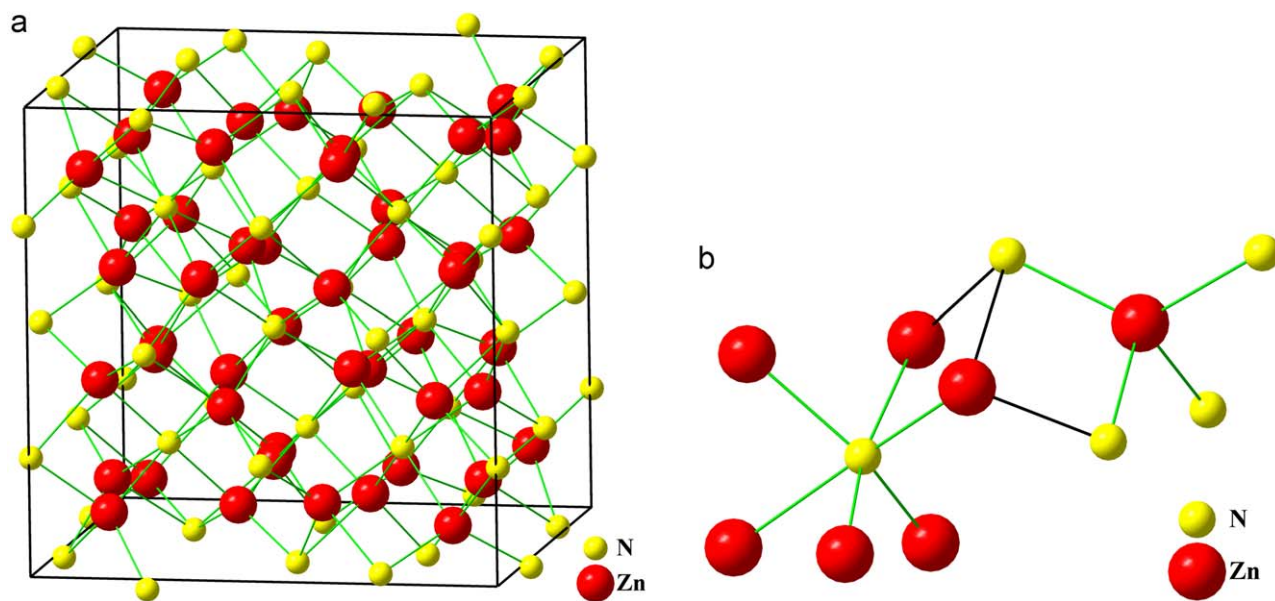


Fig. 3. (a) The schematic views of  $\text{Zn}_3\text{N}_2$ . The unit cells are outlined. (b) The details of  $\text{ZnN}_4$  tetrahedron and “ $\text{NZn}_6$  octahedron”.

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