

## Pressure effects on the transition temperature of superconducting $\text{MgC}_x\text{Ni}_3$

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The effect of hydrostatic pressure ( $P$ ) up to 17 kbar on the superconducting transition temperature ( $T_C$ ) of the newly discovered intermetallic nonoxide perovskite superconductor  $\text{MgC}_x\text{Ni}_3$  has been reported.  $T_C$  is found to increase with increasing  $P$  at a rate of  $dT_C/dP \sim 0.0134$  to  $0.0155$  K/kbar depending on the value of carbon content  $x$ . The absolute value of  $dT_C/dP$  for  $\text{MgC}_x\text{Ni}_3$  is about the same as that of intermetallic  $R\text{Ni}_2\text{B}_2\text{C}$  ( $R$  denotes rare earth) and metallic superconductors but about one order of magnitude smaller than that of the most recently and intensively studied superconductor  $\text{MgB}_2$ . However, the  $d \ln T_C/dP \sim 0.00181$  to  $0.00224$  kbar<sup>-1</sup> and the rate of change of  $T_C$  with unit cell volume ( $V$ ),  $d \ln T_C/d \ln V \sim -3.18$  to  $-2.58$  of  $\text{MgC}_x\text{Ni}_3$  are having the comparable magnitude to that of  $\text{MgB}_2$  with opposite sign. The increase of  $T_C$  with  $P$  in  $\text{MgC}_x\text{Ni}_3$  can be explained in the framework of density of states (DOS) effect.

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Soon after the discovery of a record high- $T_C$  ( $\sim 39$  K) intermetallic noncuprate superconductor  $\text{MgB}_2$ ,<sup>1</sup> a new intermetallic nonoxide superconductor  $\text{MgCNi}_3$  was found<sup>2</sup> to undergo a superconducting transition at  $T_C \sim 8$  K. Though the  $T_C$  of  $\text{MgCNi}_3$  is much lower than that of  $\text{MgB}_2$ , it still attracts a lot of attention due to at least having the following physical significance related to the present studies. (1) It has a perovskite structure as does the 30 K oxide noncuprate superconductor  $\text{Ba}_{1-x}\text{K}_x\text{BiO}_3$ .<sup>2</sup> (2) A high proportion of Ni in this compound indicates that the magnetic interactions may play a dominant role in understanding its superconductivity. (3) Its normal state NMR properties are irregular<sup>3</sup> and analogous to that observed in the exotic superconductor  $\text{Sr}_2\text{RuO}_4$ . (4) A typical isotropic  $s$ -wave superconductivity<sup>3</sup> is displayed by the nuclear spin-lattice relaxation rate ( $1/T_1$ ) with a coherence peak below  $T_C$ . (5) The change from grain boundary to core pinning by intragranular nanoparticles near  $T_C$  proposes that the arrangement of pinning sites in  $\text{MgCNi}_3$  is unique.<sup>4</sup> (6) The Hall coefficient and thermoelectric power data<sup>5,6</sup> show that the carriers in this compound are electrons in contrast to  $\text{MgB}_2$ . (7) Energy band calculations<sup>7-9</sup> demonstrate that the density of states (DOS) of the Fermi level ( $E_F$ ) is dominated by Ni  $d$  states and there is a von Hove singularity (vHS) of the DOS just below ( $< 50$  meV) the  $E_F$ .<sup>8</sup> Moreover, the photoemission and x-ray absorption studies show that the sharp vHS peak theoretically predicted near  $E_F$  is substantially suppressed which may be due to electron-electron and electron-phonon interactions.<sup>10</sup>

It is well known that the high pressure ( $P$ ) plays an important role on the  $T_C$  of the metallic and intermetallic superconductors.<sup>11-19</sup> In general,  $P$  can change the electronic structure, phonon frequencies, or electron-phonon coupling that affecting the  $T_C$ . Both positive and negative pressure derivatives,  $dT_C/dP$ , are observed in the metallic and intermetallic superconductors.<sup>11-19</sup> For example, simple

$s, p, d$ -metal superconductors<sup>16</sup> such as Sn, In, Ta, or Hg, and the intermetallic superconductor such as the recently discovered  $\text{MgB}_2$  (Refs. 11–13) have shown decreasing  $T_C$  with increasing  $P$ . However, depending on the rare earth site of the quaternary borocarbides,  $R\text{Ni}_2\text{B}_2\text{C}$  ( $R$  denotes rare earths), both an increase and decrease of  $T_C$  are observed with an increase of pressure.<sup>14,15</sup> In addition, the pressure can basically shift the Fermi level ( $E_F$ ) towards higher energies<sup>14,15</sup> and thereby provide a probe on the slope of the DOS near  $E_F$ . Moreover, it can also modify the magnetic pair breaking effect and tune the competitive phenomena between superconductivity and spin fluctuations. From our magnetic field dependent resistivity and specific-heat studies,<sup>20,21</sup> it has been suggested that the  $\text{MgCNi}_3$  is basically a typical BCS-like superconductor. In this report, we further present the pressure effects on the  $T_C$  of this exotic superconductor to testify the above-mentioned unique electronic and magnetic properties.

The details of  $\text{MgC}_x\text{Ni}_3$  sample preparation and characterization can be found in Refs. 2 and 22. With increasing the nominal carbon content,  $T_C$  was improved. Depending on the values of nominal carbon  $x$ , the samples with different  $T_C$ 's are hereafter referred as  $A$  ( $x=1.0$ ),  $B$  ( $x=1.25$ ), and  $C$  ( $x=1.5$ ). Electrical resistivity ( $\rho$ ) of  $\text{MgC}_x\text{Ni}_3$  was measured by the standard four-probe method. Thermoelectric power ( $S$ ) measurements were performed with steady state techniques. The hydrostatic pressure- ( $P$ ) dependent ac magnetic susceptibility ( $\chi_{ac}$ ) data were taken by the piston cylinder self-clamped technique.<sup>23</sup> The hydrostatic pressure environment around the sample was generated inside a Teflon cell with 3M Fluorinert FC-77 as the pressure-transmitting medium. The pressure was determined by using a Sn manometer situated near the sample in the same Teflon cell. In each instance, the original value was reproduced within experimental error after the pressure released indicating complete reversibility of the pressure effect.

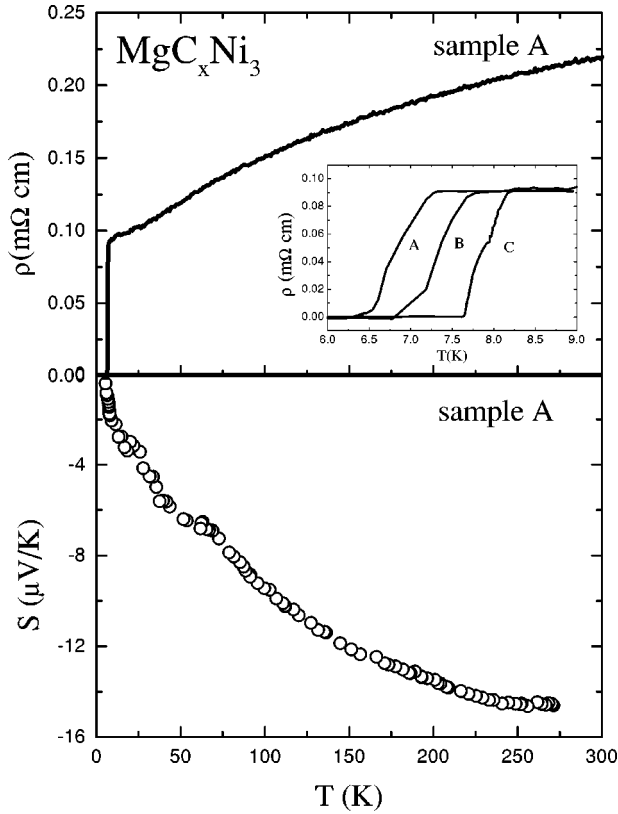


FIG. 1. Temperature ( $T$ ) variation of resistivity ( $\rho$ ) and thermoelectric power ( $S$ ) for sample A at ambient pressure. The inset shows the resistivity ( $\rho$ ) of the three samples A, B, and C near  $T_C$ .

Figure 1 shows the temperature dependence of resistivity ( $\rho$ ) and thermoelectric power ( $S$ ) for sample A. The inset of Fig. 1 displays the  $\rho$  of samples A, B, and C near  $T_C$ . The variation of  $\rho$  with temperature shows the same trend as reported in the literature<sup>2,5,6,24,25</sup> with  $T_C \sim 7-8$  K,  $\rho_{300\text{ K}}/\rho_{10\text{ K}} \sim 2.3$  and 90%-10% transition width  $\sim 0.2$  K. The different values of  $T_C$  for three studied  $\text{MgC}_x\text{Ni}_3$  samples are mainly due to the carbon stoichiometry.<sup>2,24</sup> The temperature dependence of  $S$  is negative, confirming the carriers to be electron type, which is consistent with the published results.<sup>5,6</sup> The nonlinear temperature dependence of  $S$  seems to suggest that the enhancement of electron-phonon interaction plays an important role in the superconductivity of  $\text{MgC}_x\text{Ni}_3$  as in chevre-phase compounds<sup>26</sup>  $\text{Cu}_{1.8}\text{Mo}_6\text{S}_{8-y}\text{Se}_y$  and  $\text{Cu}_{1.8}\text{Mo}_6\text{S}_{8-y}\text{Te}_y$ .

Temperature variation of ac magnetic susceptibility ( $\chi_{ac}$ ) of samples A, B, and C under pressure (0–17 kbar) is shown in Fig. 2. At ambient pressure,  $T_C$  ( $\sim 6.5$  K) of sample A is the same as that obtained from specific heat<sup>21</sup> but a little lower than that from the resistivity data (Fig. 1). The  $T_C$  (midpoint) for sample A increases from 6.56 to 6.79 K with an increase of pressure from ambient to 14.80 kbar as shown in Fig. 3, having the rate of  $dT_C/dP \sim 0.015$  K/kbar and  $d \ln T_C/dP [= (1/T_C)(dT_C/dP)] \sim 0.002$  kbar<sup>-1</sup>. The similar trend of pressure effect on  $T_C$  for samples B and C is also shown in Figs. 2 and 3. The positive values of  $dT_C/dP$  and  $d \ln T_C/dP$  for these three samples are listed in Table I. It is noted that these values of  $d \ln T_C/dP$  (Table I) for  $\text{MgC}_x\text{Ni}_3$

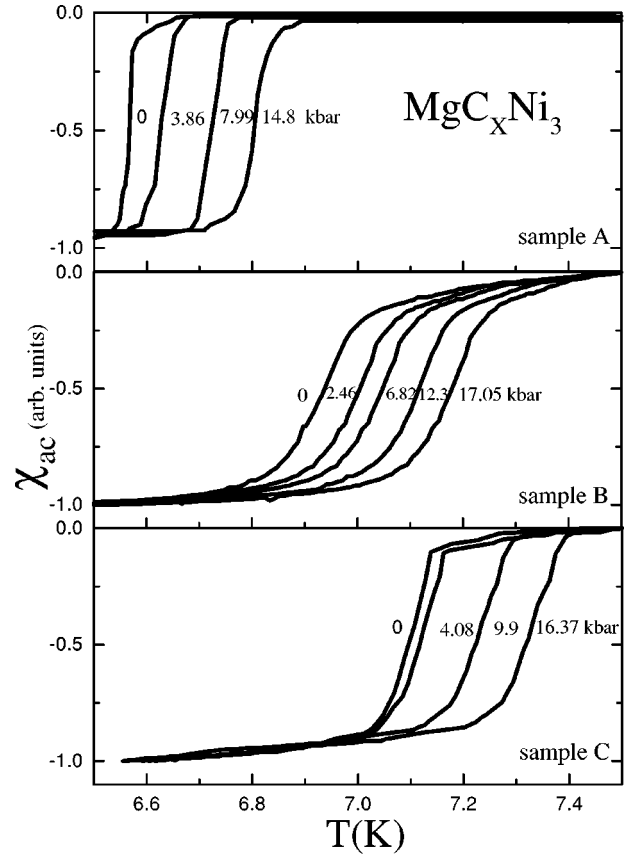


FIG. 2. Variation of ac magnetic susceptibility ( $\chi_{ac}$ ) of samples A, B, and C near  $T_C$  at various pressures ( $P$ ).

lie in the range of  $\sim 0.001-0.008$  kbar<sup>-1</sup> of conventional superconductors.<sup>27</sup>

For a clear and detailed idea of the pressure effect on the  $T_C$  of other metallic and intermetallic superconductors, some of them are also listed in Table I for comparison. The decrease and increase of  $T_C$  are observed, respectively, in metallic superconductors Ta and V with an analogous magnitude of  $dT_C/dP$  and  $d \ln T_C/dP$  as  $\text{MgC}_x\text{Ni}_3$  (Table I). It is explained by the decrease of the electron-phonon coupling constant in Ta and by the suppression of spin fluctuations as well as the increase of electron-phonon coupling in V.<sup>16-18</sup> The magnitude of positive  $dT_C/dP$  for electron-carrier  $\text{MgC}_x\text{Ni}_3$  is about the same as that of its three-dimensional analog  $\text{LuNi}_2\text{B}_2\text{C}$  superconductor, and the latter has been interpreted by an increase of the DOS with  $P$ .<sup>14</sup> However, the negative  $dT_C/dP$  and  $d \ln T_C/dP$  for hole-carrier  $\text{MgB}_2$  may be either from a decrease of the DOS (Ref. 12) or by a lattice stiffening.<sup>13</sup>

The change of  $T_C$  with the unit cell volume ( $V$ ) can be given by<sup>11,14</sup>

$$(V/T_C)(dT_C/dV) = d \ln T_C/d \ln V = -(B/T_C)(dT_C/dP), \quad (1)$$

where  $B$  is the bulk modulus of the superconductor. Using the calculated value of  $B$  for  $\text{MgC}_x\text{Ni}_3$  as 1510 kbar (Ref. 25) and taking the obtained  $dT_C/dP$  and  $T_C$  from Table I, the  $d \ln T_C/d \ln V$  values are found from Eq. (1), respec-

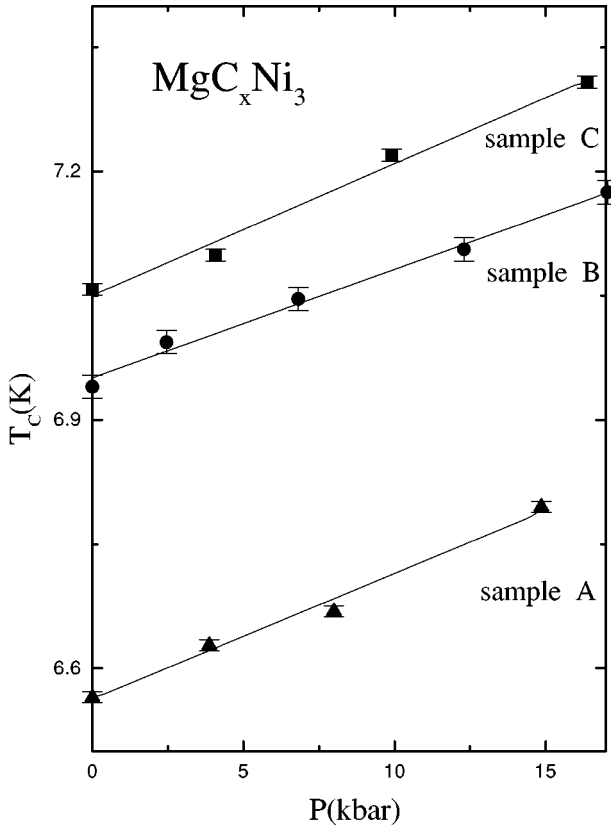


FIG. 3. Pressure ( $P$ ) dependence of superconducting transition temperature ( $T_C$ ) of samples A, B, and C.

tively, for samples A, B, and C as  $-3.18$ ,  $-2.58$ , and  $-2.76$ . These values are of the same order of magnitude in the  $\text{MgB}_2$  superconductor ( $+4.16$ ) with opposite sign.<sup>11</sup>

Since the DOS is sufficiently large in  $\text{MgC}_x\text{Ni}_3$ , produces strong electron-phonon coupling<sup>9</sup> and is supported by its  $S$  data, the  $T_C$  can be expressed by the McMillan formula<sup>28</sup> as

$$T_C = (\theta_D/1.45) \exp\{-1.04(1 + \lambda)/[1 - \mu^*(1 + 0.62\lambda)]\}, \quad (2)$$

where,  $\mu^*$  is the Coulomb pseudopotential and  $\theta_D$  is the Debye temperature.  $\lambda$  is the electron-phonon coupling constant and is given by

$$\lambda = N(E_F) \langle I^2 \rangle / M \langle \omega^2 \rangle, \quad (3)$$

where  $N(E_F)$  is the DOS at the Fermi level,  $\langle I^2 \rangle$  is the square averaged electronic matrix element for electron-phonon interaction,  $M$  is the ionic mass, and  $\langle \omega^2 \rangle$  is the square averaged phonon frequency. It appears from Eq. (2) that the change of  $\lambda$  and  $\theta_D$  by pressure will determine the sign of  $dT_C/dP$ . It is well established that the pressure induces the lattice stiffening and generally reduces the  $T_C$ .<sup>14-17</sup> However, the DOS effect can either enhance or reduce the  $T_C$  correspondingly by the increase or decrease of  $N(E_F)$  due to applied pressure.<sup>14,15</sup> The dependence of  $T_C$  on  $\theta_D$  is complicated as it appears both in the linear and exponent [being connected with  $\langle \omega^2 \rangle$  in Eq. (3)] terms of Eq. (2). Again, the change of exponent  $\lambda$  in Eq. (2) will be more effective than that of the linear term  $\theta_D$  in determining  $T_C$ .  $\theta_D$  generally increases by  $P$  amplifying the phonon frequency<sup>19</sup> as  $\langle \omega^2 \rangle = 0.5\theta_D^2$  and thus may decrease  $\lambda$  [Eq. (3)], which in turn may reduce  $T_C$  [Eq. (2)]. Therefore, the positive  $dT_C/dP$  for  $\text{MgC}_x\text{Ni}_3$  possibly originates from the increase of  $N(E_F)$  and consequently by the enhancement of electron-phonon coupling constant  $\lambda$  [Eqs. (3)] if  $\mu^*$  and  $\langle I^2 \rangle$  are less pressure dependent. In addition,  $P$  causes not only a shifting of the  $E_F$  but also a broadening of the energy bands. This broadening of energy bands may also increase  $N(E_F)$ . Most recently, Louis and Iyakutti<sup>19</sup> have successfully calculated the pressure effects on  $T_C$  of vanadium (V). Similarly, the computation of some important parameters such as  $d \ln N(E_F)/dP$  and  $d \ln \omega/dP$  of  $\text{MgC}_x\text{Ni}_3$  may be useful for quantitative analysis of our data.

Even though it is unfavorable that strong spin fluctuations exist in  $\text{MgC}_x\text{Ni}_3$ ,<sup>21</sup> the marginal or unstable spin fluctuations suppressing  $T_C$  have not been totally ruled out.<sup>8</sup> In general, pressure reduces the spin fluctuations and increases  $T_C$  because the spin fluctuations and superconductivity are mutually competitive phenomena. This may also be one of the reasons for the positive pressure effect on  $T_C$  of  $\text{MgC}_x\text{Ni}_3$ . Another considerable factor showing a positive  $dT_C/dP$  is the carbon stoichiometry in the sample. Generally, the deficiency of carbon from the optimum value decreases the  $T_C$ .<sup>2,24</sup> The nonstoichiometry of carbon (if any) may also affect the energy bands of the sample and alter the

TABLE I. The superconducting transition temperature  $T_C$  (determined from the midpoint of resistive transition for  $\text{MgC}_x\text{Ni}_3$ ) at ambient pressure,  $dT_C/dP$ , and  $d \ln T_C/dP$  for some metallic as well as intermetallic superconductors.

Sample composition	$T_C$ (K)	$dT_C/dP$ ( $10^{-2}$ K/kbar)	$d \ln T_C/dP$ ( $10^{-3}/\text{kbar}$ )	Reference
$\text{MgB}_2$	38.6	-8.0	-2.07	12
$\text{MgB}_2$	37.5	-16.0	-4.26	13
$\text{LuNi}_2\text{B}_2\text{C}$	15.9	+1.88	+1.18	14
Ta	4.3	-0.26	-0.60	16
V	5.3	+1.0	+1.88	18, 19
$\text{MgC}_x\text{Ni}_3(\text{A})$	6.9	+1.55	+2.24	This work
$\text{MgC}_x\text{Ni}_3(\text{B})$	7.4	+1.34	+1.81	This work
$\text{MgC}_x\text{Ni}_3(\text{C})$	7.9	+1.52	+1.92	This work

position of  $E_F$  compared to that expected from theoretical energy band calculations<sup>7-9</sup> for stoichiometric  $\text{MgCNi}_3$ . The present investigations for three samples with different carbon content and  $T_C$  show almost the same positive value of  $dT_C/dP$  suggesting that the carbon deficiency does not significantly affect the pressure effect of  $\text{MgC}_x\text{Ni}_3$  on  $T_C$ . However, it is noted that Kumary *et al.*<sup>25</sup> recently found a decrease of  $T_C$  up to a pressure of 17 kbar and an increase of  $T_C$  beyond this pressure using resistivity measurements. It may be possible to briefly explain these controversial results as followings. (1) The  $T_C$  determined from the resistivity (transport property) is always higher than that from susceptibility and specific heat (bulk property) measurements.<sup>2,20,21,25</sup> This may suggest that a small amount of higher  $T_C$  phase existing in the grain boundaries<sup>4</sup> superconducts through percolation effects. (2) The negative  $dT_C/dP$  observed in Ref. 25 using resistivity measurements at low pressures may be due to the reduction of grain boundary effects by pressure. Once the pressure is applied high enough

( $\sim 17$  kbar) to overcome the grain boundary effect, the bulk superconductivity dominates and the positive  $dT_C/dP$  is found to be the same with our results using susceptibility measurements.

In summary, pressure increases  $T_C$  of three intermetallic, nonoxide, and perovskite electron-type superconductors  $\text{MgC}_x\text{Ni}_3$ . The magnitude of change rate  $d \ln T_C/dP$  in  $\text{MgC}_x\text{Ni}_3$  is about the same order as that in  $\text{MgB}_2$  and  $R\text{Ni}_2\text{B}_2\text{C}$  ( $R$  denotes rare earths), which lies in the range of that of conventional superconductors. The positive value of  $dT_C/dP$  for three  $\text{MgC}_x\text{Ni}_3$  samples are almost the same and independent of various  $T_C$  resulting from different carbon stoichiometry. The present results of positive  $dT_C/dP$  of  $\text{MgC}_x\text{Ni}_3$  can be explained mainly by the increase of density of states by pressure.

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